

# A Locally Symmetric Almost Kähler Einstein Structure on the Cotangent Bundle of a Riemannian Manifold

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## ABSTRACT

We introduce a Riemannian metric of diagonal type on the cotangent bundle of a Riemannian manifold and show that  $T^*M$  with this metric is locally symmetric Einstein manifold. Also, we obtain a locally symmetric Kähler Einstein structure on the cotangent bundle of a Riemannian manifold of negative constant sectional curvature. Similar results are obtained on a tube around zero section in the cotangent bundle, in the case of a Riemannian manifold of positive constant sectional curvature.

## KEYWORDS

Almost complex structure, cotangent bundle, Kähler Einstein metric, locally symmetric Riemannian manifold.

## 1. INTRODUCTION

The differential geometry of the cotangent bundle  $T^*M$  of a Riemannian manifold  $(M, g)$  is almost similar to that of the tangent bundle  $TM$ . However, there are some differences. This is because the lifts (vertical, complete, horizontal, etc.) to  $T^*M$  cannot be defined just like in the case of  $TM$ .

In [6] V. Oproiu and D.D. Porosniuc have obtained a natural Kähler Einstein structure  $(G, J)$  of diagonal type induced on  $T^*M$  from the Riemannian metric  $g$ . The obtained Kähler structure on  $T^*M$ , depends on parameters  $u$ , which is a smooth function depending on the energy density  $\tau$  on  $T^*M$ .

In this paper, we obtain a Kähler structure on  $T^*M$ , which depends on parameters  $\alpha, \beta, u$ , where  $\alpha, \beta \in \mathbb{R}$ . The vertical distribution  $VT^*M$  and the horizontal distribution  $HT^*M$  are orthogonal to each other but the dot products induced on them from  $G$  are not isomorphic (isometric).

After that, we obtain that  $G$  is Hermitian with respect to  $J$  and it follows that the fundamental 2-form  $\Omega$ , associated to the almost Hermitian structure  $(G, J)$  is the fundamental form defining the usual symplectic structure on  $T^*M$ , hence it is closed.

From the integrability condition for  $J$  it follows that the base manifold  $M$  must have constant sectional curvature  $c$  and the parameter  $u$  must be constant.

If the constant sectional  $c$  is negative then we obtain a locally symmetric Kähler Einstein structure defined on the whole  $T^*M$ . If the constant sectional curvature  $c$  of  $M$  is positive then we get a similar structure defined on a tube around zero section in  $T^*M$ .

The manifolds, tensor fields and geometric objects we consider in this paper, are assumed to be differentiable of class  $C^\infty$  (i.e., smooth). We use the computations in local coordinates but many results from this paper may be expressed in an invariant form. The well-known summation convention is used throughout this paper, the range for the indices  $h, i, j, k, l, r, s$  being always  $\{1, \dots, n\}$  (see [5], [7]). The module of smooth vector fields on  $T^*M$  shall be denoted by  $\Gamma(T^*M)$ .

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### 1. Preliminary

Let  $(T^*M, \pi, M)$  be the cotangent bundle, where  $M$  is an  $n$ -dimensional smooth manifold. If  $(U, \varphi)$  is a local chart on  $M$  and  $(x^i)$  are the coordinates of a point  $p \in M$ ,  $p \in \varphi^{-1}(x) \in U$ , then a point  $u \in \pi^{-1}(U)$ ,  $\pi(u) = p$  has the coordinates  $(x^i, p_i)$ ,  $(i = 1, \dots, n)$ . The natural basis of the module  $\Gamma(T^*M)$  is given by  $(\partial_i = \frac{\partial}{\partial x^i}, \partial^r = \frac{\partial}{\partial p_r})$ .

An  $M$ -tensor field of type  $(r, s)$  on  $T^*M$  is defined by set of  $n^{r+s}$  components (functions depending on  $x^i$  and  $p_i$ ), with  $r$  upper indices and  $s$  lower indices, assigned to induced local charts  $(\pi^{-1}(U), \Phi)$  on  $T^*M$ , such that the local coordinate change rule is that of the local coordinate components of a tensor field of type  $(r, s)$  on the base manifold  $M$ . A usual tensor field of type  $(r, s)$  on  $M$  may be thought of as an  $M$ -tensor field of type  $(r, s)$  on  $T^*M$ . If the considered tensor field on  $M$  is covariant only, the corresponding  $M$ -tensor field on  $T^*M$  may be identified with the induced (pullback by  $\pi$ ) tensor field on  $T^*M$ .

Some useful  $M$ -tensor fields on  $T^*M$  may be obtained as follows. Let  $u, v: [0, \infty] \rightarrow \mathbb{R}$  be a smooth functions and let  $\|p\|^2 = g_{\pi(p)}^{-1}(p, p)$  be the square of the norm of the cotangent vector  $p \in \pi^{-1}(U)$ . The components  $g_{ij}(\pi(p))$ ,  $p_i$ ,  $u(\|p\|^2)p_i p_j$  define  $M$ -tensor fields of type  $(0, 2)$ ,  $(0, 1)$ ,  $(0, 2)$  on  $T^*M$ , respectively. Similarly, the components  $g^{kl}(\pi(p))$ ,  $p^i = p_h g^{hi}$ ,  $v(\|p\|^2)p^k p^l$

define  $M$ -tensor fields of type  $(2, 0)$ ,  $(1, 0)$ ,  $(2, 0)$  on  $T^*M$ , respectively. Of course, all the components considered above are in the induced local chart  $(\pi^{-1}(U), \Phi)$ .

The Levi Civita connection  $\nabla$  of  $g$  defines a direct sum decomposition

$$TT^*M = VT^*M \oplus HT^*M, \quad (1)$$

of the tangent bundle to  $T^*M$  into vertical distribution  $VT^*M = \ker \pi_*$  and the horizontal distribution  $HT^*M$ .

If  $(\pi^{-1}(U), \Phi) = (\pi^{-1}(U), x^1, \dots, x^n, p_1, \dots, p_n)$

is a local chart on  $T^*M$ , induced from the local chart  $(U, \varphi) = (U, x^1, \dots, x^n)$ , the local vector fields  $\partial^1 = \frac{\partial}{\partial p_1}, \dots, \partial^n = \frac{\partial}{\partial p_n}$  on  $\pi^{-1}(U)$  defines a local frame for  $VT^*M$  over  $\pi^{-1}(U)$  and the local vector fields  $\delta_i = \frac{\delta}{\delta x^i}, \dots, \delta_n = \frac{\delta}{\delta x^n}$  define a local frame for  $HT^*M$  over  $\pi^{-1}(U)$ , where  $\delta_i = \partial_i + p_h \Gamma_{ik}^h \partial^k$ , and  $\Gamma_{ik}^h(\pi(p))$  are the Christoffel symbols of  $g$ .

The set of vector fields  $(\partial^1, \dots, \partial^n, \delta_1, \dots, \delta_n)$

defines a local frame on  $T^*M$ , adapted to the direct sum decomposition (1).

We consider

$$\begin{aligned} \tau &= \frac{1}{2} \|p\|^2 = \frac{1}{2} g_{\pi(p)}^{-1}(p, p) \\ &= \frac{1}{2} g^{ik}(x) p_i p_k, \quad (2) \\ &p \in \pi^{-1}(U) \end{aligned}$$

the energy density defined by  $g$  in the cotangent vector  $p$ . We have for all  $p \in T^*M$ .

From now on, we shall work in a fixed local chart  $(U, \varphi)$  on  $M$  and in the induced local chart  $(\pi^{-1}(U), \Phi)$  on  $T^*M$ .

### 3. An almost Kähler structure on $T^*M$

Consider a real valued smooth function  $u$  defined on  $[0, \infty) \subset \mathbb{R}$  and real constants  $\alpha$  and  $\beta$ . We define the following symmetric  $M$ -tensor field of type  $(0, 2)$  on  $T^*M$  having the components

$$G_{ij}(p) = \frac{1}{\beta} g_{ij}(\pi(p)) + \frac{u(\tau)}{\alpha\beta} p_i p_j. \quad (3)$$

It follows easily that the matrix  $(G_{ij})$  is positive definite if and only if  $\alpha, \beta > 0$ ,  $\alpha + 2\tau u > 0$ . The inverse of this matrix has the entries

$$H^{kl}(p) = \beta g^{kl}(\pi(p)) + v(\tau) p^k p^l, \quad (4)$$

where

$$v = -\frac{u\beta}{\alpha + 2\tau u}. \quad (5)$$

The components  $H^{kl}$  define symmetric  $M$ -tensor field of type  $(2, 0)$  on  $T^*M$ . It follows:

**Remark.** If the matrix  $(G_{ij})$  is positive definite, then matrix  $(H^{kl})$  is positive definite, too.

Using the  $M$ -tensor fields defined by  $G_{ij}$  and  $H^{kl}$ , the following Riemannian metric may be considered on  $T^*M$ :

$$G = G_{ij} dx^i dx^j + H^{ij} \delta p_i \delta p_j, \quad (6)$$

where  $\delta p_i = dp_i - p_h \Gamma_{ik}^h dx^k$  is the the covariant differential of  $p_i$  with respect to the Levi Civita connection  $\nabla$  of  $g$ . Equivalently, we have

$$G(\delta_i, \delta_j) = G_{ij}, \quad G(\partial^i, \partial^j) = H^{ij}, \\ G(\delta_i, \partial^j) = G(\partial^i, \delta_j) = 0.$$

Note that  $HT^*M$  and  $VT^*M$  are orthogonal to each other with respect to  $G$ , but the Riemannian metrics induced from  $G$  on  $HT^*M$ ,  $VT^*M$  are not the same, so the considered metric  $G$  on  $T^*M$  is not a metric of Sasaki type. Note also that the system of 1-forms  $(dx^1, \dots, dx^n, \delta p_1, \dots, \delta p_n)$  defines a local frame on  $T^*T^*M$ , dual to the local frame  $(\delta_1, \dots, \delta_n, \partial^1, \dots, \partial^n)$

adapted to the direct sum decomposition (1).

Next, an almost complex structure  $J$  is defined on  $T^*M$  by the same  $M$ -tensor fields  $G_{ij}$ ,  $H^{kl}$ , expressed in the adapted local frame by

$$J(\delta_i) = G_{ik} \partial^k, \quad J(\partial^i) = -H^{ik} \delta_k \quad (7)$$

From the property of the  $M$ -tensor field  $H^{kl}$  to be defined by the inverse of the matrix defined by the components of the  $M$ -tensor field  $G_{ij}$ , it follows easily that  $J$  is an almost complex structure on  $T^*M$ .

**Theorem 1.**  $(T^*M, G, J)$  is an almost Kähler manifold.

*Proof:* Since the matrix  $(H^{kl})$  is the inverse of the matrix  $(G_{ij})$ , it follows easily that

$$G(J\delta_i, J\delta_j) = G(\delta_i, \delta_j), \quad G(J\partial^i, J\partial^j) = G(\partial^i, \partial^j), \\ G(J\delta_i, J\partial^j) = G(\delta_i, \partial^j) = 0.$$

Hence

$$G(JX, JY) = G(X, Y), \quad \forall X, Y \in \Gamma(T^*M).$$

Thus  $(T^*M, G, J)$  is an almost Hermitian manifold.

The fundamental 2-form associated with this almost Hermitian structure is  $\Omega$ , defined by

$$\Omega(X, Y) = G(X, JY), \quad \forall X, Y \in \Gamma(T^*M).$$

By a straightforward computation we get

$$\Omega(\delta_i, \delta_j) = 0, \quad \Omega(\partial^i, \partial^j) = 0, \quad \Omega(\partial^i, \delta_j) = \delta_j^i.$$

Hence

$$\Omega = \delta p_i \wedge dx^i = dp_i \wedge dx^i, \quad (8)$$

due to the symmetry of  $p_h \Gamma_{ij}^h$ . It follows that  $\Omega$  does coincide with the fundamental 2-form defining the usual symplectic structure on  $T^*M$ . Of course, we have  $d\Omega = 0$ , i.e.,  $\Omega$  is closed. Therefore,  $(T^*M, G, J)$  is an almost Kähler manifold.  $\square$

#### 4.A Kähler structure on $T^*M$

We shall study the integrability of the almost complex structure defined by  $J$  on  $T^*M$ . To do this, we need the following well-known formulas for the brackets of the vector fields  $\partial^i, \delta_j$ ,  $i = 1, \dots, n$

$$[\partial^i, \partial^j] = 0, \quad [\partial^i, \delta_j] = \Gamma_{jk}^i \partial^k, \quad [\delta_i, \delta_j] = p_h R_{kij}^h \partial^k, \quad (9)$$

where  $R_{kij}^h(\tau(p))$  are the local coordinate components of the curvature tensor field of  $\nabla$  on  $M$ . Of course, the components  $R_{kij}^h$  define  $M$ -tensor fields of type (1,3) on  $T^*M$ .

**Lemma 1.** The Nijenhuis tensor field of the almost complex structure  $J$  on  $T^*M$  is given by

$$\left\{ \begin{array}{l} N(\delta_i, \delta_j) = -\left\{ \frac{u}{\alpha\beta^2} (\delta_i^h g_{jk} - \delta_j^h g_{ik}) + R_{kij}^h \right\} p_h \partial^k, \\ N(\delta_i, \partial^j) = -H^{kl} H^{jr} \left\{ \frac{u}{\alpha\beta^2} (\delta_i^h g_{rl} - \delta_r^h g_{il}) + R_{lir}^h \right\} p_h \delta_k, \\ N(\partial^i, \partial^j) = -H^{ir} H^{jl} \left\{ \frac{u}{\alpha\beta^2} (\delta_i^h g_{rk} - \delta_r^h g_{ik}) + R_{klr}^h \right\} p_h \partial^k. \end{array} \right. \quad (10)$$

*Proof.* Recall that the Nijenhuis tensor field  $N$  defined by  $J$  is given  $N(X, Y) = [JX, JY] - J[JX, Y] - J[X, JY] - [X, Y]$ ,  $\forall X, Y \in \Gamma(T^*M)$ .

Then, we have  $\delta_k \tau = 0$ ,  $\partial^k \tau = p^k$  and  $\nabla_i G_{jk} = 0$ ,  $\nabla_i H^{jk} = 0$ , where

$$\left\{ \begin{array}{l} \nabla_i G_{jk} = \delta_i G_{jk} - \Gamma_{ij}^l G_{lk} - \Gamma_{ik}^l G_{lj}, \\ \nabla_i H^{jk} = \delta_i H^{jk} + \Gamma_{il}^j H^{lk} + \Gamma_{il}^k H^{lj}. \end{array} \right.$$

The above expressions for the components of  $N$  can be obtained by a quite long, straightforward computation.  $\square$

**Theorem 2.** Let  $M$  be a connected Riemannian

manifold and  $\dim M \geq 3$ , then the almost complex structure  $J$  on  $T^*M$  is integrable if and only if the manifold  $M$  has constant sectional curvature  $c$  and the function  $u$  is given by

$$u = -c\alpha\beta^2. \quad (11)$$

*Proof.* From the condition  $N = 0$ , one obtains

$$\left\{ \frac{u}{\alpha\beta^2} (\delta_i^h g_{jk} - \delta_j^h g_{ik}) + R_{kij}^h \right\} p_h = 0.$$

Differentiating with respect to  $p_r$ , taking  $p_h = 0 \forall h \in \{1, \dots, n\}$ , it follows that the curvature tensor field of  $\nabla$  has the expression

$$R_{kij}^r = -\frac{u(0)}{\alpha\beta^2} (\delta_i^r g_{jk} - \delta_j^r g_{ik}).$$

Using the Schur theorem it follows that  $(M, g)$  has the constant sectional curvature  $c = -\frac{u(0)}{\alpha\beta^2}$ . Then, we obtain the expression (11) of  $u$ .

Conversely, if  $(M, g)$  has constant curvature  $c$  and  $u$  is given by (11), it follows in a straightforward way that  $N = 0$ .  $\square$

It should be noted that the function  $u$  must fulfill the condition

$$\alpha + 2\tau u = \alpha(1 - 2c\beta^2\tau) > 0, \quad \alpha, \beta > 0. \quad (12)$$

Therefore, the theorem can be stated as follows:

**Theorem 3.** Let  $M$  be a Riemannian manifold with constant sectional curvature  $c$  and the function  $u = -c\alpha\beta^2$ , then

- i) If  $c < 0$  then  $(T^*M, G, J)$  is a Kähler manifold.
- ii) If  $c > 0$  then  $(T_\alpha^*M, G, J)$  is a Kähler manifold, where  $T_\alpha^*M$  is the tube around zero section in  $T^*M$  defined by the condition  $0 \leq \|p\|^2 < \frac{1}{c\beta^2}$ .
- iii) If  $c = 0$  then  $(M, g)$  is a flat manifold and  $G$  becomes a flat metric on  $T^*M$ .

The components of the Kähler metric  $G$  on  $T^*M$  are

$$\begin{cases} G_{ij} = \frac{1}{\beta} g_{ij} - c\beta p_i p_j, \\ H^{ij} = \beta g^{ij} + \frac{c\beta^3}{1 - 2c\beta^2\tau} p^i p^j. \end{cases} \quad (13)$$

## 5. A Kähler Einstein structure on $T^*M$

In this section, we shall study the property of the Kähler manifold  $(T^*M, G, J)$  to be Einstein.

The Levi Civita connection  $\bar{\nabla}$  of the Riemannian manifold  $(T^*M, G)$  is determined by the conditions

$$\bar{\nabla}G = 0, \quad T = 0,$$

where  $T$  is its torsion tensor field. The explicit expression of this connection is obtained from the formula

$$\begin{aligned} 2G(\bar{\nabla}_X Y, Z) &= X(G(Y, Z)) + Y(G(X, Z)) \\ &\quad - Z(G(X, Y)) + G([X, Y], Z) - G([X, Z], Y) \\ &\quad - G([Y, Z], X); \end{aligned}$$

$$\forall X, Y, Z \in \Gamma(T^*M).$$

The final result can be stated as follows:

**Lemma 2.** The Levi Civita connection  $\bar{\nabla}$  of  $G$  has the following expression in the local adapted frame  $(\delta_1, \dots, \delta_n, \partial^1, \dots, \partial^n)$ :

$$\begin{cases} \bar{\nabla}_{\partial^i} \partial^j = Q_h^{ij} \delta^h, \quad \bar{\nabla}_{\delta_i} \partial^j = -\Gamma_{ih}^j \partial^h + P_i^{hj} \delta_h, \\ \bar{\nabla}_{\partial^i} \delta_j = P_j^{hi} \delta_h, \quad \bar{\nabla}_{\delta_i} \delta_j = -\Gamma_{ij}^h \delta_h + S_{hij} \partial^h, \end{cases} \quad (14)$$

where  $Q_h^{ij}, P_j^{hi}, S_{hij}$  are  $M$ -tensor fields on  $T^*M$ , defined by

$$\begin{cases} Q_h^{ij} = \frac{1}{2} G_{hk} (\partial^i H^{jk} + \partial^j H^{ik} - \partial^k H^{ij}), \\ P_j^{hi} = \frac{1}{2} H^{hk} (\partial^i G_{jk} - H^{il} R_{ljk}^r p_r), \\ S_{hij} = -\frac{1}{2} G_{hk} \partial^k G_{ij} + \frac{1}{2} R_{hij}^r p_r. \end{cases} \quad (15)$$

After replacing the expressions of the involved  $M$ -tensor fields (equation (12)), we obtain

$$Q_h^{ij} = c\beta H^{ij} p_h, \quad P_j^{hi} = -c\beta H^{hi} p_j, \quad S_{hij} = c\beta G_{hj} p_i. \quad (16)$$

The curvature tensor field  $K$  of the connection  $\bar{\nabla}$  is obtained from the well-known formula

$$\begin{aligned} K(X, Y)Z &= \bar{\nabla}_X \bar{\nabla}_Y Z - \bar{\nabla}_Y \bar{\nabla}_X Z - \bar{\nabla}_{[X, Y]} Z, \\ \forall X, Y, Z &\in \Gamma(T^*M). \end{aligned}$$

The components of curvature tensor field  $K$  with respect to the adapted local frame  $(\delta_1, \dots, \delta_n, \partial^1, \dots, \partial^n)$  are obtained easily:

$$\begin{cases} K(\delta_i, \delta_j)\delta_k = c\beta(\delta_i^h G_{jk} - \delta_j^h G_{ik})\delta_h, \\ K(\delta_i, \delta_j)\partial^k = c\beta(\delta_j^k G_{hi} - \delta_i^k G_{hj})\partial^h, \\ K(\partial^i, \partial^j)\delta_k = c\beta(\delta_k^j H^{hi} - \delta_k^i H^{hj})\delta_h, \\ K(\partial^i, \partial^j)\partial^k = c\beta(\delta_h^i H^{jk} - \delta_h^j H^{ik})\partial^h, \\ K(\partial^i, \delta_j)\delta_k = c\beta\delta_j^i G_{hk}\partial^h, \\ K(\partial^i, \delta_j)\partial^k = -c\beta\delta_j^i H^{hk}\delta_h. \end{cases} \quad (17)$$

The Ricci tensor field  $Ric$  of  $\bar{\nabla}$  is defined by the formula:

$$Ric(Y, Z) = trace(X \rightarrow K(X, Y)Z),$$

$$\forall X, Y, Z \in \Gamma(T^*M).$$

It follows

$$\begin{cases} Ric(\delta_i, \delta_j) = cn\beta G_{ij}, \\ Ric(\partial^i, \partial^j) = cn\beta H^{ij}, \\ Ric(\partial^i, \delta_j) = Ric(\delta_j, \partial^i) = 0. \end{cases}$$

Thus

$$Ric = cn\beta G. \quad (18)$$

By using expression (17), we have computed the covariant derivatives of curvature tensor field  $K$  in the local adapted frame  $(\delta_i, \partial^i)$  with respect to the connection  $\bar{\nabla}$  and we obtained in all twelve cases the result is zero, i.e.,  $\bar{\nabla}K = 0$ .

Hence, we may state our main result.

**Theorem 5.** *Assume that the Riemannian manifold  $(M, g)$  has constant sectional curvature  $c$ , condition (14) are fulfilled and the components of the metric  $G$  are given by (20), then*

1. *If  $c < 0$  then  $(T^*M, G, J)$  is a locally symmetric Kähler Einstein manifold.*
2. *If  $c > 0$  then  $(T_\alpha^*M, G, J)$  is a locally symmetric Kähler Einstein manifold.*

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